

Classical Mechanics

Phys105A, Winter 2007

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Euler-Lagrange Equation

Let $y(x)$ be the path that minimizes/maximizes the integral

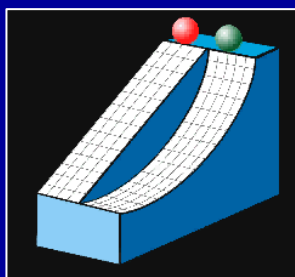
$$S = \int_{x_1}^{x_2} f[y(x), y'(x), x] dx$$

The Euler-Lagrange equation tells us that S is extremal, or *stationary*, when $y(x)$ obeys

$$\frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial y'} = 0$$

Fastest Path

Given two points P, Q in \mathbb{R}^2 in a gravitational field $g\mathbf{e}_x$, what is the fastest path $y(x)$?



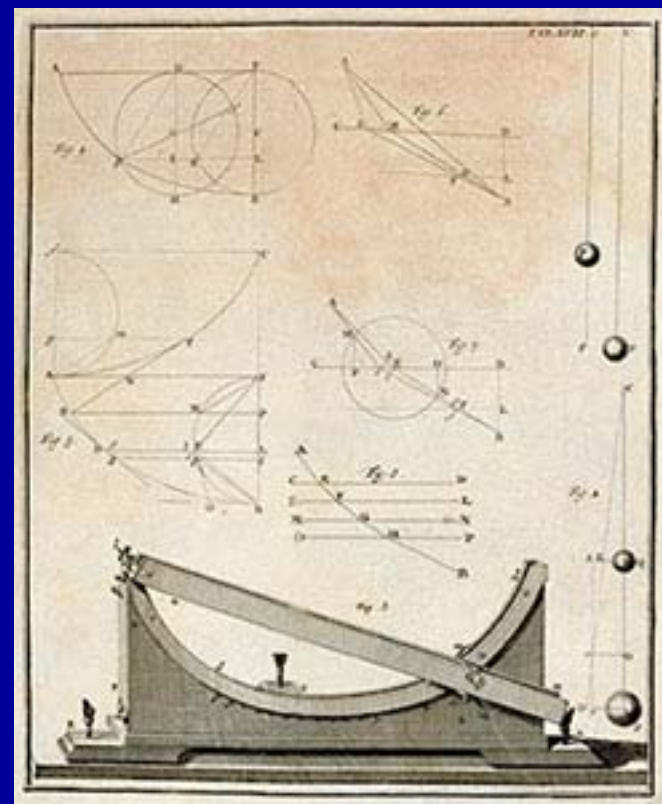
Find $y(x)$ that minimizes

$$\frac{1}{\sqrt{2g}} \int_{x_1}^{x_2} \frac{\sqrt{1+y'^2}}{\sqrt{x}} dx$$

The answer to this classic *brachistochrone problem* is that $y(x)$ is (part) of a cycloid

$$x(\theta) = a(1 - \cos\theta)$$

$$y(\theta) = a(\theta - \cos\theta)$$



Brachistochrone Problem

For the shortest path we look for the extremal values of the integral $\int f(x) dx$ with $f(x) = \sqrt{(1+y'^2)}/\sqrt{x}$.

Euler-Lagrange says $\partial f/\partial y - d(\partial f/\partial y')/dx = 0$, hence $\partial f/\partial y'$ is constant (in x).

$$\frac{\partial f}{\partial y'} = \frac{1}{\sqrt{x}} \frac{y'}{\sqrt{1+y'^2}} = c \quad \text{gives the equation} \quad y' = \sqrt{\frac{x}{2a-x}}$$

$$y(x_2) = \int_{x=0}^{x_2} \sqrt{\frac{x}{2a-x}} dx \quad \text{has solution} \quad \begin{aligned} x &= a(1 - \cos \theta) \\ y &= a(\theta - \sin \theta) \end{aligned}$$

E-L for Several Variables

For an arbitrary number of variables $q_1(t), \dots, q_n(t)$, with t the independent variable, what are the extremals of:

$$S = \int_{t_1}^{t_2} L[q_1, \dot{q}_1, \dots, q_n, \dot{q}_n, t] dt ?$$

The Euler-Lagrange equation tells us (again) that the dependent variables have to obey for all j :

$$\frac{\partial L}{\partial q_j} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} = 0$$

Notation

Note the following shorthand conventions:

The derivative of y with respect to x can be denoted using the apostrophe: $dy/dx = y'$.

Derivatives with respect to time are often denoted with a dot above the coordinate. Hence

$$y' = \frac{dy}{dx} \quad \text{and} \quad \dot{y} = \frac{dy}{dt}$$

The “ ∂ versus d ” notation in $\partial L/\partial q - d(\partial L/\partial \dot{q})/dt = 0$ should remind you that t is the independent variable and that d/dt concerns all variables in L (unlike the partial derivatives such as $\partial L/\partial q$).

Chapter 7



Lagrange's Equations

Langrangian

For an unconstrained system in 3 dimensions subject to the a conservative field with potential energy $U=U(\mathbf{r})$ and kinetic energy $T= \frac{1}{2}m(x'^2+y'^2+z'^2)$ we have the **Langrangian (or Lagrange function)**: $\mathcal{L} = T-U$, in this case $\mathcal{L}(x,y,z,x',y',z') = \frac{1}{2}m(x'^2+y'^2+z'^2) - U(x,y,z)$.

Note that $\partial\mathcal{L}/\partial x = -\partial U/\partial x = F_x$ and $\partial\mathcal{L}/\partial x' = mx' = p_x$, hence $\partial\mathcal{L}/\partial x - d(\partial\mathcal{L}/\partial x')/dt = 0$; similarly for y and z .

These are the three Lagrange equations, which follow from Newton's second law $\frac{\partial\mathcal{L}}{\partial q} = \frac{d}{dt} \frac{\partial\mathcal{L}}{\partial \dot{q}}$ for $q(t) = x(t)$, $y(t)$, and $z(t)$.